

QUASILINEAR SPIN VOLTAGE PROFILES IN SPIN THERMOELECTRICS

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Contribution of conduction electrons to Spin Seebeck effect

- ferromagnet with paramagnetic impurities → inelastic spin-flip scattering
- diffusion contribution at low temperatures

Questions:

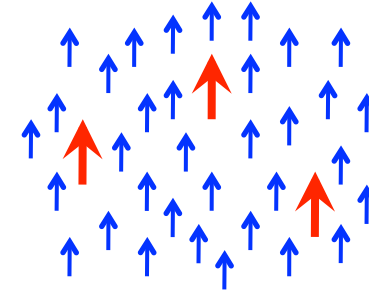
- quasilinear variation of spin voltage possible?
- relation of spin voltage to spin thermopower?

Model

Itinerant ferromagnet: $\varepsilon_{p\sigma} = \varepsilon_p - \sigma h$
 with **paramagnetic impurities**

Exchange interaction

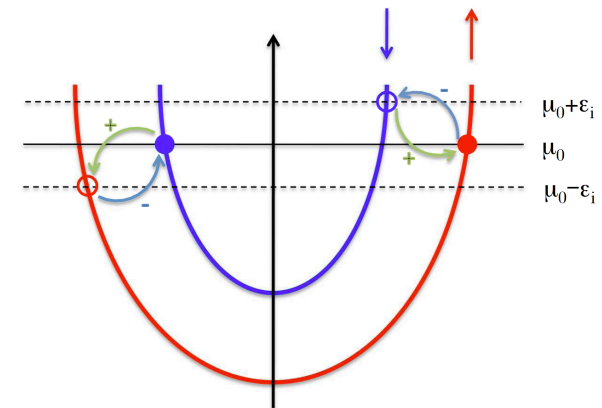
→ splitting of impurities S_z -levels by energy ε_i



Method: Kinetic equation → derive continuity equations for currents \mathbf{J}_σ and $\mathbf{J}_{Q\sigma}$
 for weak ferromagnets at low temperatures

Relaxation mechanisms:

- Inelastic spin-flip scattering between electrons and impurity spins (rate $\Gamma_\sigma^{\text{sf}}$)
asymmetric for \uparrow and \downarrow spins at Fermi energy
(finite $\Gamma_-^{\text{sf}} = N_\uparrow \Gamma_\uparrow^{\text{sf}} - N_\downarrow \Gamma_\downarrow^{\text{sf}}$)
- Direct spin relaxation of electrons (rate $\Gamma_\sigma^{\text{dir}}$)
 (e.g. spin-orbit interaction)
- Direct spin relaxation of impurities (rate Γ_i^{dir})
 (e.g. spin-orbit interaction + phonons)



Spin-electric effects ($\nabla T = 0$)

For linear spin voltage μ_- need:

- Conservation law
- Chemical potential μ_+ as sink/source on right hand side

Continuity equations:

$$\nabla \cdot \mathbf{J} = -\frac{\sigma_+}{e^2} \nabla^2 \mu_+ - \frac{\sigma_-}{e^2} \nabla^2 \mu_- = 0 \quad \text{Conservation law}$$

$$\nabla \cdot \mathbf{J}_{\text{spin}} = -\frac{\sigma_-}{e^2} \nabla^2 \mu_+ - \frac{\sigma_+}{e^2} \nabla^2 \mu_- = -2\alpha \Gamma_-^{\text{sf}} \mu_+ - 2(\alpha \Gamma_+^{\text{sf}} + \Gamma_+^{\text{dir}}) \mu_-$$

direct spin relaxation of impurities

spin asymmetry of inelastic spin-flip scattering

$$\alpha = \frac{\Gamma_i^{\text{dir}}}{\Gamma_i^{\text{sf}} + \Gamma_i^{\text{dir}}}$$

$$\Gamma_-^{\text{sf}} = N_{\uparrow} \Gamma_{\uparrow}^{\text{sf}} - N_{\downarrow} \Gamma_{\downarrow}^{\text{sf}}$$

Linear charge voltage μ_+ \rightarrow linear spin voltage μ_-

$$\mu_- \simeq -\mu_+ \frac{\alpha \Gamma_-^{\text{sf}}}{\alpha \Gamma_+^{\text{sf}} + \Gamma_+^{\text{dir}}}$$

Spin thermoelectrics – linear μ_-

Need to determine voltages μ_+ and μ_- and temperatures T_+ and T_-

→ consider continuity equation for currents: \mathbf{J} , \mathbf{J}_{spin} , \mathbf{J}_Q , $\mathbf{J}_{Q\text{spin}}$

μ_+ and T_+ appear on right hand side of continuity equations due to finite Γ_-^{sf}

→ decoupling of continuity equations impossible

Linear μ_- when **additional conservation law** (besides $\nabla \cdot \mathbf{J} = 0$) exists

- no direct spin relaxation for electrons and impurities $\nabla \cdot \mathbf{J}_{\text{spin}} = \nabla \cdot \mathbf{J}_Q = 0$
- no direct spin relaxation for impurities $\nabla \cdot \mathbf{J}_Q = 0$
- no direct spin relaxation for electrons $\nabla \cdot (\epsilon_i \mathbf{J}_{\text{spin}} + 2\mathbf{J}_Q) = 0$

Spin thermoelectrics – quasilinear μ_-

In general: direct spin relaxation exists for both electrons and impurities

→ currents \mathbf{J}_{spin} , \mathbf{J}_Q , $\mathbf{J}_{Q\text{spin}}$ decay on length scales:

$$\frac{1}{\ell_1^2} = |e|^2 \frac{\Gamma_+^{\text{sf}} + \Gamma_+^{\text{dir}}}{\sigma_+}, \quad \frac{1}{\ell_2^2} = \frac{1}{\ell_1^2} \frac{\alpha \Gamma_+^{\text{sf}} + \Gamma_+^{\text{dir}}}{\Gamma_+^{\text{sf}} + \Gamma_+^{\text{dir}}},$$

$$\frac{1}{\ell_3^2} = \frac{1}{\ell_1^2} \frac{\alpha \epsilon_i \Gamma_-^{\text{sf}} \Gamma_+^{\text{dir}}}{2(\Gamma_+^{\text{sf}} + \Gamma_+^{\text{dir}})(\alpha \Gamma_+^{\text{sf}} + \Gamma_+^{\text{dir}})} \left(\frac{e}{2\gamma} S_+ - \frac{\Gamma_-^{\text{sf}'}}{\Gamma_-^{\text{sf}}} \right)$$

- ℓ_1 ordinary spin diffusion length
- ℓ_2 of same order than ℓ_1 but diverges for $\Gamma_\sigma^{\text{dir}} = \Gamma_i^{\text{dir}} = 0$
- ℓ_3 generically much longer than ℓ_1 and diverges for $\Gamma_\sigma^{\text{dir}} = 0$ or for $\Gamma_i^{\text{dir}} = 0$

$$\frac{\ell_3}{\ell_1} \text{ is of order } \begin{cases} \frac{\epsilon_f}{\epsilon_i} \sqrt{\frac{\Gamma_+^{\text{dir}}}{\Gamma_+^{\text{sf}}}} \gg 1 & \text{for } \Gamma_+^{\text{dir}} \gg \Gamma_+^{\text{sf}} \\ \frac{\epsilon_f}{\epsilon_i} \sqrt{\frac{\Gamma_+^{\text{sf}}}{\Gamma_+^{\text{dir}}}} \gg 1 & \text{for } \Gamma_+^{\text{sf}} \gg \Gamma_+^{\text{dir}} \end{cases}$$

Spin voltage

Ferromagnet with constant temperature gradient:



No direct spin relaxation: $\mu_- = -e (S_- T_+ + S_+ T_-)$

→ conventional expression

General case: direct spin relaxation exists for both electrons and impurities

$$\mu_- = \alpha \frac{\frac{e}{2} S_+ \Gamma_-^{\text{sf}} - \gamma \Gamma_-^{\text{sf}'}}{\alpha \Gamma_+^{\text{sf}} + \Gamma_+^{\text{dir}}} T_+ - \gamma \frac{\alpha \Gamma_+^{\text{sf}'} + \Gamma_+^{\text{dir}'}}{\alpha \Gamma_+^{\text{sf}} + \Gamma_+^{\text{dir}}} T_-$$

→ spin thermopower S_- does not contribute

Outlook

Quasilinear contribution of conduction electrons to spin Seebeck effect in ferromagnets with paramagnetic impurities due to asymmetry of inelastic spin-flip scattering

So far: diffusion contribution to spin voltage at low temperatures

Future goal: extend theory to higher temperatures

→ need to consider:

- Electron-electron interactions
- Scattering off phonons (phonon drag)
- Scattering off magnons (magnon drag)